

Multiband Superconductivity in Heavy Fermion Compound CePt₃Si without Inversion Symmetry: An NMR Study on a High-Quality Single Crystal

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We report on novel superconducting characteristics of the heavy fermion (HF) superconductor CePt₃Si without inversion symmetry through ¹⁹⁵Pt-NMR study on a single crystal with $T_c = 0.46$ K that is lower than $T_c \sim 0.75$ K for polycrystals. We show that the intrinsic superconducting characteristics inherent to CePt₃Si can be understood in terms of the unconventional strong-coupling state with a line-node gap below $T_c = 0.46$ K. The mystery about the sample dependence of T_c is explained by the fact that more or less polycrystals and single crystals inevitably contain some disordered domains, which exhibit a conventional BCS *s*-wave superconductivity (SC) below 0.8 K. In contrast, the Néel temperature $T_N \sim 2.2$ K is present regardless of the quality of samples, revealing that the Fermi surface responsible for SC differ from that for the antiferromagnetic order. These unusual characteristics of CePt₃Si can be also described by a multiband model; in the homogeneous domains, the coherent HF bands are responsible for the unconventional SC, whereas in the disordered domains the conduction bands existing commonly in LaPt₃Si may be responsible for the conventional *s*-wave SC. We remark that some impurity scatterings in the disordered domains break up the *4f*-electrons-derived coherent bands but not others. In this context, the small peak in $1/T_1$ just below T_c reported before [Yogi *et al.* (2004)] is not due to a two-component order parameter composed of spin-singlet and spin-triplet Cooper pairing states, but due to the contamination of the disorder domains which are in the *s*-wave SC state.

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1. Introduction

For a superconducting order parameter, it is generally believed that time reversal invariance is a necessary condition for the formation of spin-singlet Cooper pairing, while an inversion center is additionally required for the formation of spin-triplet pairing.¹⁾ In superconductors that lack inversion symmetry, the relationship between spatial symmetry and the Cooper-pair spin state may be broken, which will make the parity of superconducting state mixed between even and odd parities due to the antisymmetric spin-orbit coupling (ASOC),²⁻⁸⁾ thereby leading to a two-component order parameter composed of spin-singlet and spin-triplet Cooper pairing states. The first heavy fermion (HF) superconductivity (SC) without inversion symmetry was discovered in CePt₃Si.⁹⁾ After that, the pressure-induced superconductivity of this family was also observed for a ferromagnet UIr¹⁰⁾ and antiferromagnets CeRhSi₃,¹¹⁾ CeIrSi₃,¹²⁾ and CeCoGe₃.¹³⁾ An intriguing feature of these systems is that the splitting of the Fermi surfaces by ASOC is much larger than the superconducting gap.^{5,14)} Interestingly, in the SC state of CeIrSi₃ and CeRhSi₃, it has been reported that the upper critical fields (H_{c2} 's) exceed 30 T along the *c*-axis, which are extremely high as compared to a

Pauli limiting field.¹⁵⁻¹⁷⁾ This kind of behavior has never been observed in the other HF superconductors such as CeCoIn₅, CeIrIn₅, and CeCu₂Si₂, which are in a spin-singlet Cooper pairing regime. In the noncentrosymmetric superconductors where the correlation between electrons may be not so significant, however, the characteristic feature for conventional *s*-wave spin singlet superconducting state has been observed in many compounds, for example, Y₂C₃,^{18,19)} Ir₂Ga₉,^{20,21)} LaPt₃Si,²²⁾ LaIrSi₃,²³⁾ Li₂Pd₃B,²⁴⁾ and so on, whereas the parity mixing state has been argued only in the related compound Li₂Pt₃B.^{25,26)}

The CePt₃Si exhibits superconductivity at $T_c = 0.75$ K in an antiferromagnetically ordered state below a Néel temperature of $T_N = 2.2$ K, as reported by Bauer *et al.*^{9,27)} Neutron-scattering measurement probed an antiferromagnetic (AFM) structure with a wave vector $\mathbf{Q} = (0, 0, 1/2)$ and a magnetic moment of $0.16 \mu_B$ lying in the *ab*-plane of the tetragonal lattice.²⁸⁾ Uniform coexistence of the AFM order and SC has been microscopically evidenced by NMR^{29,30)} and μ SR.³¹⁾ Although CePt₃Si has attracted considerable attention as the first superconductor without inversion symmetry,³⁻⁸⁾ the experimental results are still contradictory. For example, the London penetration depth,³²⁾ thermal conductivity,³³⁾ and NMR²⁹⁾ revealed the line node in the SC gap function well below T_c , whereas the small peak that was observed in $1/T_1$ just below T_c provided evidence for the inclusion of an *s*-

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wave component without nodes.²⁹⁾ However, the evidence of parity mixing due to the lack of inversion symmetry has not been obtained in this compound so far, suggesting that the theoretical models for CePt₃Si put forth to date⁶⁾ seem to be inadequate because of the unknown complicated multiband effect. The T_c value, however, still remains a mystery;^{34–38)} Takeuchi and coworkers reported that for a high-quality single crystal of CePt₃Si, specific-heat measurements revealed a bulk SC at $T_c = 0.46$ K,³⁹⁾ which is remarkably lower than $T_c = 0.75$ K for a polycrystal.⁹⁾ Despite the lower T_c of the sample, its quality is guaranteed by the extremely small value of a residual γ term, i.e., the T -linear coefficient of the electronic specific heat well below T_c is smaller than that for polycrystals.³⁹⁾ In this high-quality single crystal, however, the resistivity drops to zero below 0.75 K, similar to the case of polycrystals. There still remains an underlying mystery, namely, the sample dependence of T_c prevents us from identifying any intrinsic SC properties of CePt₃Si. In particular, it is also crucial to reveal the origin of the small peak in $1/T_1$ just below T_c in the previous sample with $T_c \sim 0.75$ K.²⁹⁾

On the basis of ¹⁹⁵Pt-NMR studies on a high-quality single crystal, we report in this paper that the genuine SC characteristics inherent to CePt₃Si can be understood in terms of an unconventional strong-coupling state with a line–node gap below $T_c = 0.46$ K; further, the mystery about the sample dependence of T_c is explained by the fact that the disordered domains with $T_c = 0.8$ K are inevitably contained even in a single crystal. Then, we propose a possible multiband model to account for the disparate SC characteristics and the universal presence of the AFM order in CePt₃Si. This model allows us to consider that SC and the AFM order occur on different pieces of Fermi surfaces, but the $4f$ -electrons-derived coherent HF bands collapse partially due to the disorders inevitably included in the samples. In this context, the small peak in $1/T_1$ just below T_c reported in the previous paper²⁹⁾ is not due to a two-component order parameter composed of spin-singlet and spin-triplet Cooper pairing states, but due to the contamination of the disorder domains that are in the s -wave SC state. Furthermore, we compare the results on CePt₃Si with that of another noncentrosymmetric HF compound CeIrSi₃. As a result, we remark that multiband and single band pictures may be applicable for understanding the phase diagrams of AFM and SC in CePt₃Si and CeIrSi₃, respectively, leading to an indication that the higher SC transition tends to emerge for a single band rather than for multibands among the various HF compounds.

2. Experimental

A high-quality single crystal (#4) of CePt₃Si was grown by the Bridgman method in a Mo crucible, as described elsewhere.³⁹⁾ A very small piece of the crystal aligned along the (001) tetragonal plane was cut out from the ingot of a large single crystal by using the x-ray Raue method. The quality of this sample was guaranteed by a sharp jump in the specific heat at $T_c = 0.46$ K, and by the small value of the residual γ term, which was less than 34 mJ/(K²·mol) well below T_c . The resistivity of the sample began to decrease below 0.8 K; at 0.75 K, the resistivity became zero, as reported in literature.³⁹⁾ For NMR measurements, this crystal

Table I. CePt₃Si samples used in our NMR studies. The previous NMR studies by Yogi *et al.*^{29,30)} were performed mainly for the single crystal #2.⁴³⁾

Sample	Crystal	T_c^{SHb} (K)	T_c^{Rc} (K)	Reference
#1	poly. (coarse powder)	~0.75	~0.75	9
#2	single (coarse powder)	0.4–0.8	~0.75	29, 30, 43
#3a	poly. ^{a)} (coarse powder)	0.4–0.8	~0.75	29, 30, 43
#3b	poly. ^{a)} (well-ground powder)	~0.75	~0.75	29, 30, 43
#4	single (coarse powder)	~0.46	~0.75	39

a) ²⁹Si-enriched sample

b) T_c^{SH} : estimated from a specific heat jump

c) T_c^{R} : estimated from a drop in the resistivity

(#4) was moderately crushed into coarse powder; the particles were approximately more than a few hundred micrometers in diameter and were oriented with the magnetic field parallel to the c -axis.

We examined the sample dependence of SC and the magnetic properties of CePt₃Si via NMR studies. The samples used in the NMR studies are listed in Table I along with their T_c values estimated from the jump in the specific heat and the drop in the resistivity. The previous NMR investigations by Yogi *et al.*^{29,30)} were carried out using a single crystal (#2), as in literature.⁴³⁾ The specific-heat measurement of this sample revealed a broad increase below 0.75 K, and a small peak at 0.46 K (see Fig. 4); further, the resistivity dropped to zero below 0.75 K. It should be noted that the polycrystal sample (#1) used in the first report by Bauer *et al.*⁹⁾ exhibited a broad peak in the specific heat and zero resistance around $T_c = 0.75$ K, but did not reveal any anomalies around 0.46 K. In order to examine the possible disorder effects arising from lattice distortion, the NMR study was also carried out by crushing coarse powder of ²⁹Si-enriched polycrystal (#3a) into well-ground powder (#3b). The average diameters of #3a and #3b were approximately a few hundred micrometers and less than 20 μm , respectively. In particular, the sample #3b is anticipated to include many local lattice distortions or crystal defects in the shorter length scale than the diameter of the particles due to the thoroughly crushing process over several hours. It should be noted that the temperature (T) dependence of the specific heat of the samples (#3a) and (#3b) are different, as will be shown in Fig. 6. Thus, unexpected sample dependence of the SC property is one of the main issues in this study, as discussed later.

3. Results and Analyses

3.1 ¹⁹⁵Pt-NMR spectra

In order to characterize the samples used for the NMR studies, we present the field-swept ¹⁹⁵Pt-NMR spectra for the single crystal (#4) and other crystals in Fig. 1. Two inequivalent crystallographic Pt sites denoted by Pt(1) and Pt(2)²⁹⁾ are distinguished by the difference in their Knight shifts. The NMR spectral widths for the single crystals #4 and #2 and the polycrystal #3a are narrower than those for the polycrystal #1 and the well-ground powder sample #3b; from these spectral widths we can observe that the single crystal #4 has better quality on a microscopic level. In contrast, polycrystals #1 and #3b are affected by some

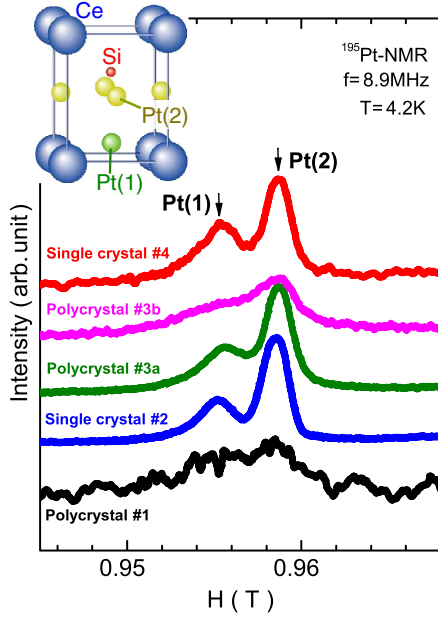


Fig. 1. (Color online) ^{195}Pt -NMR spectra at 8.9 MHz and 4.2 K for the single crystal #4 and other crystals. The inset indicates the atomic configuration in a unit cell of CePt_3Si without inversion symmetry. Two inequivalent crystallographic Pt sites in a unit cell — denoted by Pt(1) and Pt(2) — are distinguished by the difference in their Knight shifts.²⁹ The NMR spectral widths for the single crystal #4, the single crystal #2, and the polycrystal #3a are narrower than those for the polycrystal #1⁹⁾ and the well-ground powder sample #3b, indicating that the samples #1 and #3b include some kind of disorders.

disorders such as impurities, crystal defects, local-lattice distortions, and so on. It should be noted that the T dependence of the ^{195}Pt Knight shift for the present single crystal #4 coincides with that of #2.^{29,30)}

3.2 Nuclear spin–lattice relaxation rate

The measurements of the nuclear spin–lattice relaxation rate $1/T_1$ allow us to characterize each sample. When an electronic state is homogeneous over a sample, the recovery curve of ^{195}Pt nuclear magnetization ($I = 1/2$) is generally determined by a simple exponential function given by,

$$m(t) \equiv \frac{M(\infty) - M(t)}{M(\infty)} = \exp\left(-\frac{t}{T_1}\right),$$

where $M(\infty)$ and $M(t)$ are the nuclear magnetization at the thermal equilibrium condition and the nuclear magnetization at a time t after the saturation pulse, respectively. However, as indicated in Fig. 2, $m(t)$ cannot be fitted by a single exponential function for all the crystals. Unexpectedly, even in the high-quality sample #4, the $m(t)$ cannot be fitted by a single exponential function; however, it can be fitted by a multiexponential function over the entire temperature range. Such behavior was also observed in the recovery curves at both the Pt(1) and Si sites, indicating the presence of some inhomogeneity in the electronic states over the sample, which are spatially distributed over a macroscopic scale. It is noteworthy that the fraction of long components in T_1 is larger for the polycrystal #1 than for the single crystal #4. Furthermore, the fraction of the long components of the well-ground powder sample #3b becomes larger than that for the coarse powder #3a. These facts reveal that the longer

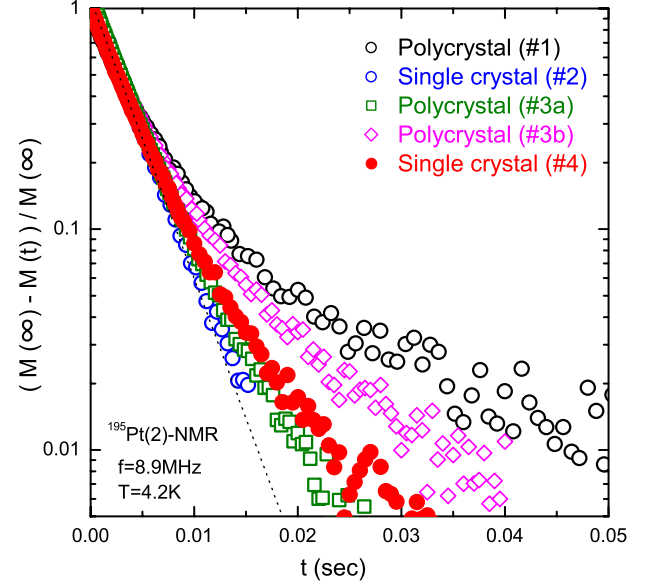


Fig. 2. (Color online) Recovery curves of nuclear magnetization $m(t)$ (see text) at the Pt(2) site for the single crystal #4 and other samples at 4.2 K. It should be noted that all $m(t)$ curves cannot be fitted by a single exponential function (dotted line), but by a multiexponential function. A similar behavior was also observed at the Pt(1) and Si sites, suggesting that the electronic states of all the samples were more or less distributed over each sample, depending on a fraction of the disordered domains into which defects were inevitably or intentionally introduced. The long component (T_{1L}) becomes more prominent for the polycrystal #1 and the well-ground powder sample #3b, which represents the large volume fractions of the disordered domains in their crystals.

components arise from the domains affected by some disorders such as impurities, crystal defects, local-lattice distortions, and so on.

The recovery curve of the single crystal #4 is tentatively fitted by assuming two components in T_1 , i.e., short and long components denoted by T_{1S} and T_{1L} , respectively, as follows:

$$m(t) = A_S \exp\left(-\frac{t}{T_{1S}}\right) + A_L \exp\left(-\frac{t}{T_{1L}}\right),$$

where A_S and A_L are the respective fractions of the short and long components, and $A_S + A_L = 1$. It should be noted that $A_S \approx 0.7$ and $A_L \approx 0.3$ for the single crystal #4 correspond to the respective volume fractions of the homogeneous and disordered domains in the sample; hence, it can be stated that these fractions depend on the samples.

3.3 Intrinsic SC properties inherent to CePt_3Si : Short T_1 component of single crystal #4

First we note a short component T_{1S} that is more dominant for the high-quality single crystal (#4), which represents an intrinsic property inherent to CePt_3Si . Figure 3 shows the T dependence of $1/T_{1S}T$ at $H \sim 0.96$ T along the c -axis. It decreases below $T_N = 2.2$ K and drops steeply below $T_c(H) \approx 0.35$ K at $H \sim 0.96$ T for the single crystal with $T_c = 0.46$ K at $H = 0$.³⁹⁾ Interestingly, in the SC state, the $1/T_{1S}$ shows a T^3 -like dependence without a coherence peak just below T_c , followed by a $T_1T = \text{const.}$ -like behavior well below T_c . These facts are evidence of the unconventional superconducting nature of the intrinsic domain of CePt_3Si . By assuming a line–node gap model with $\Delta = \Delta_0 \cos \theta$ and

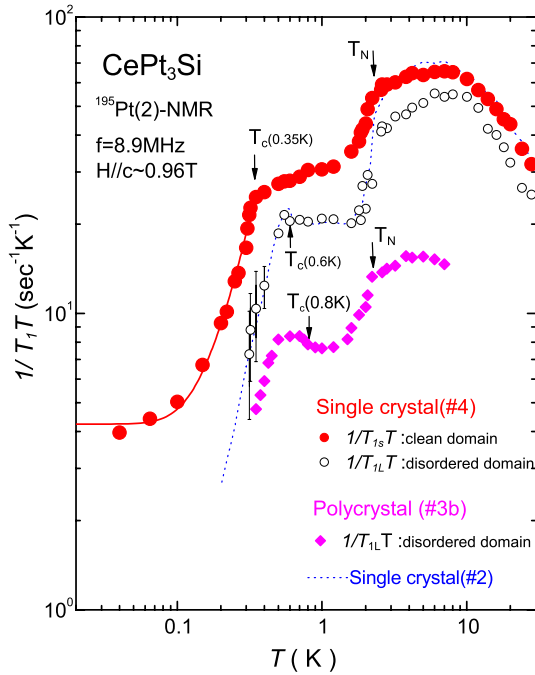


Fig. 3. (Color online) T dependence of $1/T_{1S}T$ and $1/T_{1L}T$ at the Pt(2) site of the single crystal #4. $1/T_{1S}T$ derived from the homogeneous domains evidences the unconventional SC state with the line–node gap below $T_c(H) = 0.35$ K at $H_c = 0.96$ T along the c -axis. In contrast, $1/T_{1L}T$ derived from the disordered domains decreases with a tiny peak just below $T_c(H) \sim 0.6$ K, which is the same as the T_c for the polycrystal at $H \sim 0.96$ T, similar to the previous result for the single crystal #2 (dotted line) obtained by Yogi *et al.*²⁹

the residual density of states (RDOS) at the Fermi level $N_{\text{res}}(E_F)$, $1/T_{1S}T$ in the SC state is well reproduced by

$$\frac{1/T_{1S}T}{1/T_{1S}T_c} = \frac{2}{k_B T} \int \left(\frac{N_s(E)}{N_0} \right)^2 f(E)[1 - f(E)] dE,$$

where $N_s(E)/N_0 = E/\sqrt{E^2 - \Delta^2}$; N_0 is the DOS at E_F in the normal state and $f(E)$ is the Fermi distribution function. As shown by the solid line in Fig. 3, the experimental result is well fitted by assuming $2\Delta_0/k_B T_c \approx 6$ and $N_{\text{res}}/N_0 \approx 0.41$, suggesting that strong-coupling SC emerges with line–node gap. It is expected that the large RDOS in the SC state is caused due to the impurity effect and/or the Volovik effect where the RDOS is induced by a supercurrent in the vortex state.⁴⁰ Since the contribution of the former is estimated to be less than 10% of the DOS at T_c from specific-heat measurement,³⁹ the contribution of the latter may be dominant in the present case.

3.4 Disordered domain in the single crystal: Long T_1 component of single crystal #4

The T dependence of the long component $1/T_{1L}T$ for the single crystal #4 is different from that of the short T_1 component. As shown in Fig. 3, the $1/T_{1L}T$ decreases below $T_N = 2.2$ K, and there is a marked decrease in $1/T_{1L}T$ with a tiny peak just below $T_c(H) \sim 0.6$ K that is the value in the field of $H \sim 0.96$ T for the polycrystal with $T_c(H = 0) = 0.75$ K.⁹ This T variation of $1/T_{1L}T$ resembles the result reported previously for the single crystal #2,²⁹ as shown by the dotted line in Fig. 3. Although we have claimed that the $1/T_1$ in the single crystal #2 has a peak just below T_c in the

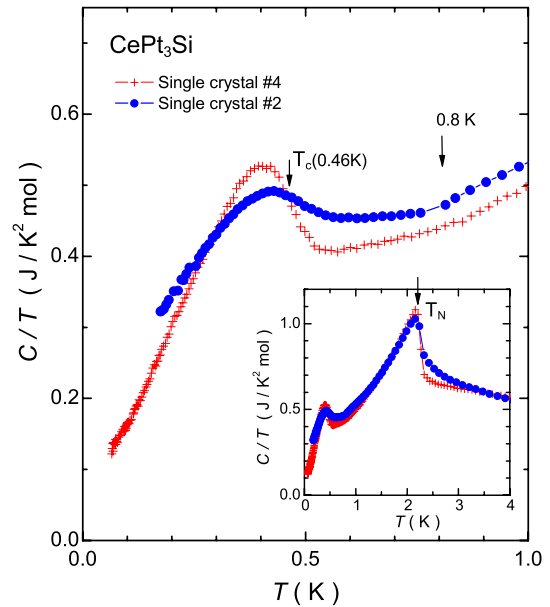


Fig. 4. (Color online) The specific heat in the single crystal #2 used in the previous NMR study²⁹ along with that in the present crystal #4 (cited from ref. 39). It exhibits the gradual increase of C/T between 0.4 and 0.8 K upon cooling and broader peak at 0.46 K than in the present single crystal #4, suggesting that the sample #2 actually includes large inhomogeneity in the SC property due to the presence of the disordered domains.

previous paper,²⁹ we should note that the previous result represents the $1/T_1$ data of the disordered domains accidentally contained. In fact, the previous single crystal #2 includes large inhomogeneity in the SC property, which was suggested by the specific heat measurement that exhibits the gradual increase of C/T between 0.4 and 0.8 K upon cooling and broader peak at 0.46 K than in the present single crystal #4,³⁹ as shown in Fig. 4. If we extract $1/T_{1S}T$ in the sample #2 by assuming two T_1 components, its T -dependence resembles the present results on the sample #4. Eventually, as will be discussed about the well-ground powder sample #3b in the next section, the origin of the small peak of $1/T_1$ for the single crystal #2²⁹ is relevant with the coherence peak just below T_c (~ 0.8 K) in the disorder-rich domains. Nevertheless, it is puzzling that $T_c = 0.75$ K for the disordered domains is higher than that for the homogeneous domains, namely, $T_c = 0.46$ K, despite the fact that disorders generally decrease T_c in unconventional HF superconductors.

3.5 SC characteristics of disorder-rich domains: Long T_1 of a well-ground powder sample #3b

In order to shed light on the SC characteristics of disordered domains, we examined a well-ground powder sample #3b into which disorders were intentionally introduced. The presence of disorder in this sample was actually suggested by the facts of the significant line broadening of the NMR spectrum (see Fig. 1) and a large fraction of long T_1 components in the recovery curve (see Fig. 2). The $m(t)$ of this sample (#3b) comprises the multi T_1 components due to the distribution of the disordered domains. A predominant component of the long T_{1S} is extracted from the $m'(t)$ that is obtained by subtracting the contribution of short T_{1S}

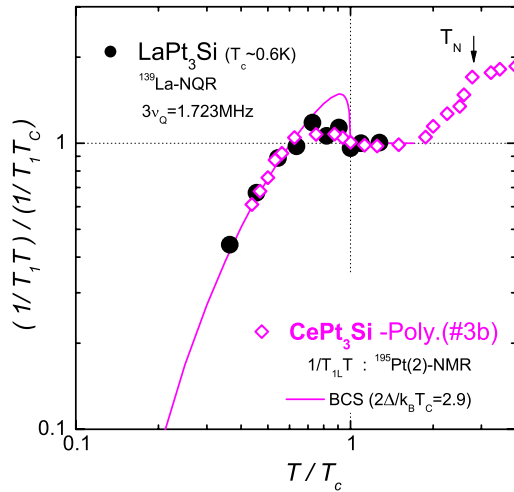


Fig. 5. (Color online) $1/T_{1L}T$ normalized by the value at T_c for sample #3b, whose electronic states are dominated by disordered domains. A distinct coherence peak in $1/T_{1L}T$ just below $T_c \approx 0.8$ K and the exponential decrease in $1/T_{1L}T$ well below T_c reveal that the disordered domains are in the BCS s -wave SC regime, resembling the s -wave superconductor LaPt_3Si with $T_c = 0.6$ K.²²⁾ The solid curve indicates the result calculated by means of an isotropic s -wave model with $2\Delta_0/k_B T_c = 2.9$.

component, which represents the SC characteristics of the disorder-rich domains.

In the normal state, the $1/T_{1L}T$ of the disordered domain of the sample #3b decreases below $T_N \sim 2.2$ K and stays a constant just above 0.8 K, as shown in Fig. 3. Remarkably, a distinct coherence peak appears just below $T_c \approx 0.8$ K, followed by an exponential-like decrease in $1/T_{1L}T$ well below T_c . Unexpectedly, this result resembles the result of $1/T_1T$ of the isostructural compound LaPt_3Si with $T_c = 0.6$ K. Figure 5 shows the $1/T_{1L}T$ normalized by the value at T_c for the disordered domains of well-ground powder #3b in comparison with that of LaPt_3Si investigated by ^{135}La -NQR at $3\nu_Q(\pm 7/2 \leftrightarrow \pm 5/2) = 1.723$ MHz.²²⁾ As shown by the solid curve in the figure, these results are reproduced by an isotropic gap model with $2\Delta_0/k_B T_c \approx 2.9$ in a weak coupling regime of a BCS-type s -wave superconductor mediated by the electron–phonon interaction.

Here, we note that the values of T_c for the disordered domains of CePt_3Si , namely, 0.8 K, is equal to the SC onset temperature in resistivity and specific-heat measurements for most CePt_3Si crystals.^{9,34–39)} It is suggested that the drop of the resistivity in CePt_3Si below 0.8 K takes place in the disordered domains that are inevitably contained. It is noteworthy that the disorders introduced by grinding the crystals of CePt_3Si into fine powder increase the domains with higher SC transition, namely, $T_c \approx 0.8$ K. Indeed, the increase of bulk T_c is also corroborated by the specific heat measurement for the well-ground powder sample #3b. As shown in Fig. 6, a broad jump in C/T was observed at 0.8 K in the case of sample #3b after grinding; the coarse-powdered polycrystal sample #3a before grinding exhibited a small jump in C/T at 0.46 K in addition to a gradual increase in C/T below 0.8 K. In the case of sample #3b, which contains more disordered domains, the residual γ term increases well below T_c . In particular, the T dependence of

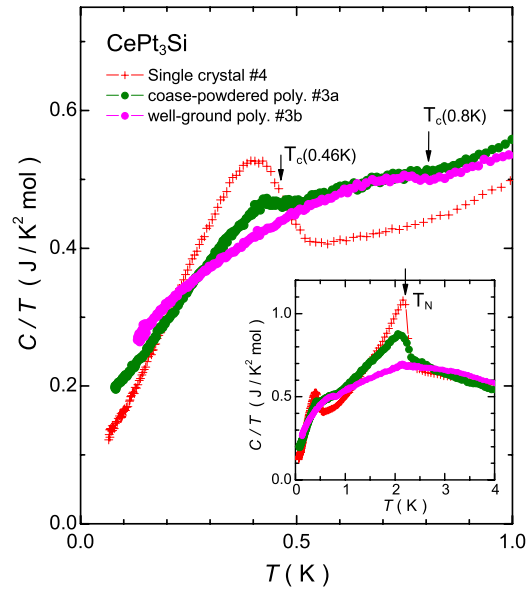


Fig. 6. (Color online) T dependence of specific heat divided by temperature for the well-ground powder sample #3b and the coarse powder sample #3a, along with the result for the single crystal #4 of CePt_3Si . The small jump in C/T at 0.46 K in the coarse powder #3a disappears in the well-ground powder sample #3b. The T dependence of C/T for sample #3b is similar to that for the polycrystal #1 first reported by Bauer *et al.*⁹⁾ This result is consistent with the T_1 result in that the fraction of T_{1L} increases in the case of the sample #3b and the polycrystal #1, both of which are more affected by disordered domains than the others are, as shown in Fig. 2.

C/T for the sample #3b is quite similar to that for the polycrystal sample #1 reported by Bauer *et al.*⁹⁾ This result is consistent with the results of the NMR spectrum (Fig. 1) and the recovery curve (Fig. 2). From these results, the SC state with high T_c (0.8 K) for the disordered domains is consistent with the characteristics of Anderson’s dirty superconductor theorem for the s -wave Cooper pairing state with an isotropic gap.⁴¹⁾ The effects of the disorders on thoroughly crushing the LaPt_3Si sample is not observed from the recovery curve and the $1/T_1$ data, which suggests that a conduction band inherent to LaPt_3Si is not affected by the disorders. This is also true for the SC state of the disordered domains of CePt_3Si . Nevertheless, it should be noted that the conventional SC state of the disordered domains of CePt_3Si coexists with the AFM order on a microscopic scale, as evidenced by the decrease in $1/T_{1L}T$ below $T_N \sim 2.2$ K, as shown in Fig. 3. In this context, we suppose that the coexistence of SC and the AFM order in the disordered domains in CePt_3Si resembles the magnetic superconductors in ternary compounds RERh_4B_4 ($\text{RE} = \text{rare earth}$), where the conventional SC derived from the conduction electron coexists with the localized $4f$ -electrons-derived magnetic order.⁴²⁾ Thus, it should be stated that the SC state of the disordered domains in CePt_3Si is not entirely the same as that of LaPt_3Si . In fact, the s -wave SC in CePt_3Si survives under the application of magnetic field of 0.96 T parallel to the c -axis, whereas that of LaPt_3Si is suppressed even by $H = 0.1$ T,³⁹⁾ revealing that the coherence length and the penetration depth of the disordered domains of CePt_3Si would be different from those of LaPt_3Si .

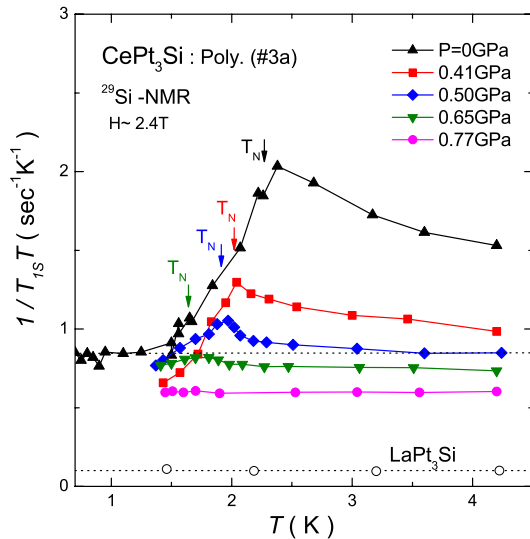


Fig. 7. (Color online) T dependence of $1/T_{1s}T$ measured by ^{29}Si -NMR for the polycrystal #3a of CePt_3Si under P . The $1/T_{1s}T$ dominated by the AFM spin fluctuations above T_N at $P = 0$ is completely suppressed for $P > 0.7$ GPa, in association with the drastic decrease of T_N determined by the peak of $1/T_{1s}T$. The DOS at E_F below T_N , which is proportional to the value of $\sqrt{1/T_{1s}T}$ for $T_c < T \ll T_N$, does not change significantly with the pressure. This result indicates that the collapse of the AFM order does not influence the Fermi surfaces that are responsible for the onset of SC.

3.6 Pressure effect on CePt_3Si : A ^{29}Si -NMR of ^{29}Si -enriched polycrystal (#3a)

Next, we address the pressure effect on the AFM order in CePt_3Si through ^{29}Si -NMR measurements of $1/T_{1s}T$ under P for the homogeneous domains of the ^{29}Si -enriched polycrystal #3a. The AFM order of CePt_3Si is completely suppressed when $P > 0.7$ GPa^{39,43,44} [see Fig. 8(a)], whereas its T_c decreases slightly around an AFM quantum critical point (QCP). As shown in Fig. 7, the $1/T_{1s}T$ enhanced by the AFM spin fluctuations above T_N is gradually reduced by applying P and is completely suppressed when $P > 0.7$ GPa, in association with the drastic decrease in T_N determined by the peak of $1/T_{1s}T$. An interesting observation is that the value of $1/T_{1s}T$ for $T_c < T \ll T_N$ is $\sim 0.8 \text{ s}^{-1} \text{ K}^{-1}$ when $P = 0$, which is comparable to $\sim 0.6 \text{ s}^{-1} \text{ K}^{-1}$ when $P = 0.77$ GPa, irrespective of the drastic suppression of the AFM order. Since the fraction of DOS at E_F below T_N is proportional to the value of $\sqrt{1/T_{1s}T}$ for $T_c < T \ll T_N$, it is revealed that the DOS at the Fermi surface, which is responsible for the onset of SC, does not change when the AFM order collapses. This result reveals that the Fermi surface relevant with the onset of SC differs from that with the AFM order.

4. Discussion

4.1 Multiband scenario for CePt_3Si

The SC properties of CePt_3Si are unusual. The SC inherent to CePt_3Si is unconventional with a line–node gap of $T_c = 0.46$ K; however, the disorders suppress this intrinsic SC, instead, the s -wave SC emerges as a result of the increase of the disordered domains, thereby resulting in a high value of T_c of almost 0.8 K. In contrast with the SC characteristics that differ according to the quality of the samples, the AFM order with $T_N \sim 2.2$ K is robust against the disorders. These results reveal that the Fermi surfaces

relevant with the onset of SC differs from that with the AFM order and that the disorders suppress the unconventional SC emerging in the $4f$ -electrons-derived HF bands inherent to CePt_3Si . These facts allow us to account for by using a multiband model, that is, there exist three characteristic bands at least. The first one is a localized f -electron band far below the Fermi level that causes the long-range AFM order below 2.2 K. It should be noted that the large reduction in $1/T_1$ points to a possible opening of the gap at the Fermi level below T_N . The second one is the f -electrons-derived coherent HF bands that are derived from the formation of a periodic lattice of Ce atoms. The third one is a weakly correlated conduction band which leads to the onset of s -wave SC in LaPt_3Si . In the homogeneous domains in samples, the coherent HF bands, which are formed through the hybridization between f electrons and conduction electrons, are responsible for the unconventional SC. Since the disorders break up the coherence of the periodicity of Ce atoms to suppress the unconventional SC inherent to the HF bands, eventually, the conduction bands in the inhomogeneous domains commonly found in LaPt_3Si may be responsible for the conventional s -wave SC possibly mediated by the electron–phonon interaction. In this model, if the sample were ideally homogeneous being free from any disorders, the conventional BCS s -wave SC taking place on the conduction bands may be unfavorable due to the pair-breaking originating from the strong electron correlation effect through an interband coupling with the HF bands. Consequently, it is likely that the decrease in resistivity below $T_c \approx 0.8$ K is due to the inevitably introduced disordered domains. Simultaneously, the specific heat shows a broad peak or a double peak between $T_c = 0.46$ K for the homogeneous domains and $T_c \approx 0.8$ K for the disordered domains,^{9,34–36,39} which depends on a fraction of the disordered domains to the homogeneous ones. This is a possible explanation to settle several underlying issues reported thus far, for instance, the sample dependence of T_c .

4.2 Comparison of the SC characteristics of CePt_3Si with that of CeIrSi_3

Here, the SC of CePt_3Si is compared with the pressure (P)-induced SC of CeIrSi_3 , since both compounds lack inversion symmetry along the c -axis. Interestingly, as illustrated in Figs. 8(a)^{39,43,44} and 8(b),¹² the P – T phase diagrams of AFM and SC of CePt_3Si and CeIrSi_3 are very different. It should be noted that the T_c for CeIrSi_3 is enhanced up to 1.6 K by virtue of the presence of strong AFM spin fluctuations around QCP.²³ In the case of CePt_3Si , T_c ($= 0.46$ K) for CePt_3Si decreases progressively with pressure, and its SC emerges under the Fermi-liquid state without any trace of AFM spin fluctuations. In order to gain insight into this contrasted SC characteristics of both compounds, we would remark a multiband effect in HF systems by systematically comparing with most of HF SC compounds reported thus far.

A number of studies on HF compounds have revealed that unconventional SC arises at or close to the QCP, where the magnetic order disappears at low temperatures as the function of P . These findings suggest that the mechanism forming Cooper pairs can be magnetic in origin. However, the natures of SC and magnetism are still unclear when SC

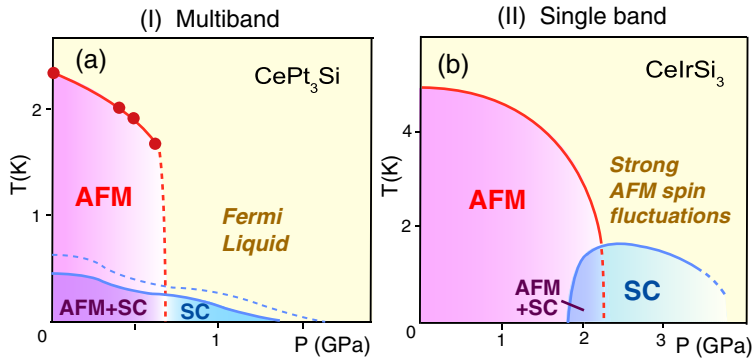


Fig. 8. (Color online) The P - T phase diagram for (a) CePt_3Si ^{39,43,44} and (b) CeIrSi_3 .¹² The SC inherent to CePt_3Si emerges under the Fermi-liquid state without any trace of AFM spin fluctuations, whereas that of CeIrSi_3 occurs under the Fermi liquid state with strong AFM spin fluctuations.

appears very close to AFM. In fact, the P -induced first-order transition from AFM to paramagnetism has been revealed in CeIn_3 ,^{45,46} CePd_2Si_2 ,⁴⁷ and CeRh_2Si_2 .^{48,49} near the boundary where SC emerges without the development of AFM spin fluctuations. For these compounds—denoted as group [I] compounds—that exhibit relatively high T_N and low T_c ,⁵⁰ a multiband model was proposed considering that the FS relevant with the AFM order differs from that with SC,⁴⁶ similar to the case of CePt_3Si . Remarkably, a different behavior was reported in the archetypal HF superconductors— CeCu_2Si_2 ^{51–54} and CeRhIn_5 ^{55–57}—with relatively low T_N and high T_c ⁵⁰ (denoted as group [II] compounds). Although an analogous behavior relevant with an AFM-QCP has been demonstrated in both compounds, it is noteworthy that the SC region associated with group [II] extends to higher pressures than that in group [I], their T_c value reaching its maximum away from the verge of the AFM order.^{56–59} The SC of group [II] compounds emerges under the Fermi-liquid state dominated by strong AFM spin fluctuations. From a comparison of group [II] with group [I], we propose that the SC state in group [II] may be described by a single band picture, namely, both AFM and SC take place on the same band. Interestingly, in high- T_c copper oxides where the single band picture is applicable, the phase diagram of AFM and SC has been recently established as the function of the carrier doping level,^{60,61} resembling that of group [II] compounds. Although the band structure of HF compounds is generally very complicated, there exist typical experimental results that indicate the validity of the single band picture where AFM and SC occur commonly. For example, (1) the dHvA experiment with CeRhIn_5 under P ⁶² revealed that the FS topology dramatically changes at the QCP where the localized- f electronic states turn into the coherent HF bands. (2) The corresponding magneto-transport measurements revealed the anomaly at the QCP on CeMIn_5 ($M = \text{Rh}$ and Co)⁶³ and also the high- T_c cuprates. (3) In the uniformly coexisting phase of AFM and SC of CeRhIn_5 under $P = 1.6$ GPa, the NQR measurement revealed that the extremely large RDOS remained at the Fermi level below T_c because of an intimate coupling between the AFM order parameter and the SC order parameter,⁵⁷ whereas the RDOS became very small in the single SC phase at $P \sim 2$ GPa. In contrast, in the case of group [I], the HF bands remaining below T_N are responsible for the unconventional SC with the line-node gap, as well as in group [II]; however, a very small residual DOS at E_F was observed well below T_c , even in the uniformly coexisting

phase of AFM and SC for CeIn_3 ,⁴⁶ UPd_2Al_3 ,⁶⁴ URu_2Si_2 ,⁶⁵ and CePt_3Si . In this context, a criterion that the single band picture is applicable would be whether a large fraction of RDOS remains well below T_c . As a result, the AFM interaction in the same band plays vital role not only for the AFM order but also for the strong coupling SC in group [II]. Once the AFM order collapses above QCP, the AFM order parameter will totally evolve into the SC order parameter, resulting in relatively high SC transition around QCP. Furthermore, we note that the tetracritical point of the AFM phase, the uniformly mixed phase of AFM+SC, the SC phase, and the PM phase is present at a certain temperature and pressure without a magnetic field in CeRhIn_5 ⁵⁷ and high- T_c copper oxides.⁶¹ In this context, we propose that these phenomena in group [II] and high- T_c copper oxides may be phenomenologically described by the $\text{SO}(5)$ theory that unifies the AFM and SC states by a symmetry principle and describes their rich phenomenology through a single low-energy effective model.⁶⁶ Therefore, on the basis of these findings, we consider that CePt_3Si and CeIrSi_3 are classified into the groups [I] and [II], respectively. Consequently, these considerations will lead us to a coherent understanding for the SC around the QCP of the AFM order in strongly correlated matter.

4.3 Possible stacking faults in the high-quality single crystals

In the CePt_3Si , no evidence of parity mixing due to the lack of inversion symmetry has been obtained experimentally so far, suggesting that the theoretical models for CePt_3Si put forth to date⁶ seem to be inadequate because of the unknown complicated multiband effect. There still remains the underlying issue regarding why the disordered domains—amounting to 30% of the volume fraction—are present even in the highest quality single crystal (#4) of CePt_3Si . Unexpectedly, these domains could be distinguished neither by X-ray spectroscopy nor by the resonance line shape in NMR. It should be noted that a such marked sample dependence of T_c has never been reported in CeIrSi_3 ¹² and CeRhSi_3 ,¹¹ and also in other HF superconductors.

One of possible explanation for the nonnegligible amount of the disordered domain in high-quality single crystal is the presence of stacking faults that are inevitably introduced in the process of crystal growth. As illustrated in Figs. 9(a) and 9(b), the stacking faults are assumed to be present along the c -axis on a microscopic scale; the intrinsic and disordered domains of the high-quality single crystal are derived from

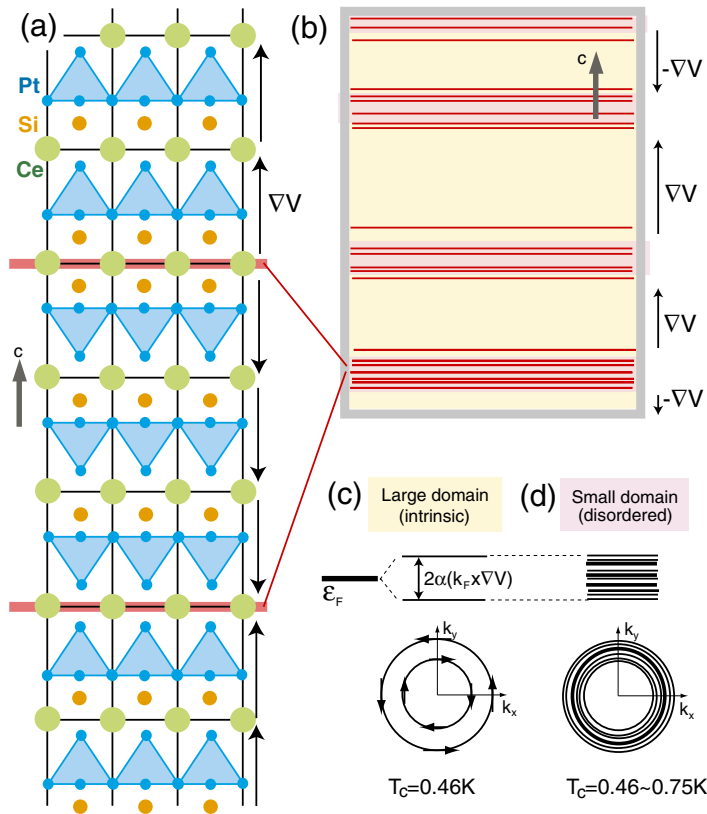


Fig. 9. (Color online) The present single crystal (#4) includes a large volume fraction of homogeneous domains with $T_c = 0.46$ K, including a small fraction of disordered domains with $T_c = 0.8$ K. (a) As one of the possible origins of the disordered domains in this sample, we assume that stacking faults are present along the c -axis on a microscopic scale, as indicated by the red lines. Here, the inversion symmetry is seemingly recovered at an interface. (b) A large-scale view of the crystal, including dense stacking faults. (c) For the homogeneous domains, the Rashba-type ASOC causes the splitting of the Fermi surfaces due to the lack of inversion symmetry. (d) The ASOC is reduced as a result of an average over interfaces for dense stacking faults with many interfaces. This may also make the disordered domains break up the coherent HF bands to suppress the unconventional SC even in the single crystal.

large twin domains with less interfaces and small twin domains with many interfaces, respectively. In this case, the inversion symmetry is expected to be seemingly recovered at the interfaces of the stacking fault, resulting in the reduction and/or distribution of the Rashba-type ASOC locally. In the homogeneous domains, an asymmetric potential gradient $\mathbf{E} = \nabla V \parallel (001)$ brings about the Rashba-type ASOC $\mathcal{H}_{\text{SO}} \sim \alpha(\mathbf{k} \times \nabla V) \cdot \sigma$.²⁾ Here, α is the coupling constant of the ASOC. This causes the splitting of the Fermi surfaces (FSs) into two pieces; the electron spins align clockwise on one FS and anticlockwise on the other FS so as to be parallel to $\mathbf{k}_F \times \nabla V$, where \mathbf{k}_F is the Fermi momentum, as illustrated in Fig. 9(c). In the disordered domains, however, the asymmetric potential gradient is locally reduced and/or distributed, which causes a possible variation in the amplitude and sign of ∇V . As illustrated in Fig. 9(d), this event is anticipated to reduce and/or distribute the Rashba-type ASOC and hence to give a degeneracy again on the clockwise and anticlockwise spin states on these FSs. This may also act as one of the scattering sources that causes quasiparticles to break up the f -electron-derived coherent HF bands as well as the disorders induced in the well-ground sample do.

5. Conclusion

¹⁹⁵Pt-NMR studies on CePt₃Si have revealed the presence of homogeneous and disordered domains even in a high-quality single crystal. Homogeneous domains inherent to CePt₃Si exhibit unconventional SC with a line-node gap below $T_c = 0.46$ K. In these domains, the coherent HF bands, which are formed through the hybridization between f -electrons and conduction electrons, are responsible for the unconventional SC. In contrast, the disordered domains reveal the conventional BCS s -wave SC with a high T_c of

0.8 K, which is analogous to $T_c = 0.6$ K for LaPt₃Si. In the disordered domains, the conduction bands existing commonly in LaPt₃Si may be responsible for the conventional s -wave SC state possibly mediated by the electron-phonon interaction. In this context, the small peak in $1/T_1$ just below T_c reported in the previous paper²⁹⁾ is not due to a two-component order parameter composed of spin-singlet and spin-triplet Cooper pairing states, but due to the contamination of the disorder domains that are in the s -wave SC state.

A ²⁹Si-NMR study under pressure has revealed that the collapse of the AFM order does not significantly change the fraction of the DOS at E_F that is responsible for the onset of SC. This result provides sound evidence that the Fermi surfaces that are responsible for the SC state differ from those responsible for the AFM order. On the basis of these results, we propose that these unusual SC and magnetic characteristics of CePt₃Si can be described by a multiband model in which it is considered that the impurity scattering is caused by the disorders and/or that the stacking faults break up the $4f$ -electrons-derived coherent HF bands, but not others. In order to gain further insight into the possible order parameter in CePt₃Si, Knight-shift measurement in the SC state using a better single crystal is highly desired.

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