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# Dynamic nuclear polarization enhanced magnetic field sensitivity and decoherence spectroscopy of an ensemble of near-surface nitrogen-vacancy centers in diamond

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We perform pulsed optically detected electron spin resonance to measure the DC magnetic field sensitivity and electronic spin coherence time  $T_2$  of an ensemble of near-surface, high-density nitrogen-vacancy centers engineered to have a narrow magnetic resonance linewidth. Combining pulsed spectroscopy with dynamic nuclear polarization, we obtain the photon-shot-noise-limited DC magnetic sensitivity of  $35 \text{ nT Hz}^{-0.5}$ . We find that  $T_2$  is controlled by instantaneous diffusion, enabling decoherence spectroscopy on residual nitrogen impurity spins in the diamond lattice and a quantitative determination of their density. The demonstrated high DC magnetic sensitivity and decoherence spectroscopy are expected to broaden the application range for two-dimensional magnetic imaging. *Published by AIP Publishing.* [<http://dx.doi.org/10.1063/1.4983350>]

Submicron scale, two-dimensional (2D) magnetic imaging has potential applications in biological and physical sciences.<sup>1,2</sup> The realization of utilizing an ensemble of nitrogen-vacancy (NV) centers in diamond is particularly attractive due to its high magnetic sensitivity at ambient conditions.<sup>3–10</sup> When NV-based sensing is carried out with continuous wave (CW) optically detected magnetic resonance (ODMR), the photon-shot-noise-limited DC magnetic field sensitivity  $\eta_{\text{sn}}$  is estimated as

$$\eta_{\text{sn}}^{(\text{cw})} = \frac{h}{g\mu_B} \frac{\delta\nu}{C\sqrt{I_0}}, \quad (1)$$

where  $h/g\mu_B = 36 \text{ } \mu\text{T/MHz}$  is the inverse of the gyromagnetic ratio of the NV electronic spins,  $I_0$  is the count rate of photons from the NV centers in a unit area ( $1 \text{ } \mu\text{m}^2$ ) under the off-resonance condition,  $\delta\nu$  is the ODMR linewidth, and  $C$  is the ODMR contrast (the ratio of the photon counts on and off resonances).<sup>11,12</sup> Equation (1) suggests that simultaneously achieving a high NV density and a narrow linewidth is desired to improve  $\eta_{\text{sn}}^{(\text{cw})}$ . In addition, the NV sensor must be located as close as possible to a magnetic specimen; it is crucial to have an NV ensemble near the diamond surface.<sup>13,14</sup>

Recently, some of the present authors have reported the successful creation of a 100-nm-thick layer of NV ensembles at a diamond surface with a density of  $10^{17} \text{ cm}^{-3}$  and a  $\delta\nu$  of  $\sim 200 \text{ kHz}$ .<sup>15</sup> This was achieved by a combination of

chemical vapor deposition (CVD) growth of a nitrogen-doped, nuclear-spin-free  $^{12}\text{C}$  (99.9%) diamond film<sup>16</sup> and subsequent helium ion implantation to introduce vacancies into the film. The detailed procedure for the NV formation and the characterization of the NV ensemble by photoluminescence (PL) spectroscopy and CW ODMR are given in Ref. 15.

In this paper, we show that, by concurrently applying pulsed ODMR and dynamic nuclear polarization (DNP) techniques to the same diamond sample, it is possible to realize an  $\eta_{\text{sn}}$  of  $35 \text{ nT Hz}^{-0.5}$ . We also examine the coherence properties of the NV ensemble to extract quantitative information on residual paramagnetic impurities in the sample. This is an example of “decoherence spectroscopy/imaging,” in which magnetic signals are detected via the change in spin coherence time  $T_2$ .<sup>17–23</sup> This method is applicable to identify magnetic signals external to the sample, providing another tool for ensemble-based 2D magnetic field imaging.

We first recap the main result of Ref. 15 by performing CW ODMR at an external magnetic field  $B_0$  of 1.5 mT. The squares ( $\square$ ) in Fig. 1 are the measured  $\delta\nu$  (top) and  $C$  (middle) together with  $\eta_{\text{sn}}^{(\text{cw})}$  estimated from Eq. (1) (bottom) as functions of the microwave power  $P_{\text{mw}}$ , demonstrating the minimum sensitivity of  $124 \text{ nT Hz}^{-0.5}$ . The measurement setup in the present work is a home-built confocal microscope combined with microwave circuitry, enabling CW and pulsed ODMR of single and ensemble NV centers. Throughout this work,  $B_0$  is applied parallel to one of four NV axes, and the  $m_S = 0 \leftrightarrow -1$  transition of the NV ensemble aligned to the field is examined, unless otherwise

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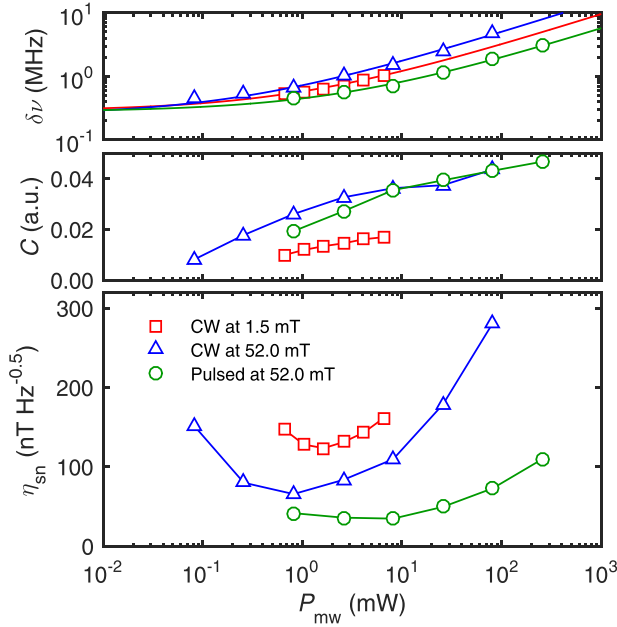


FIG. 1.  $\delta\nu$  (top),  $C$  (middle), and  $\eta_{\text{sn}}$  (bottom) as functions of  $P_{\text{mw}}$  measured at the input port of the PCB board on which the sample is mounted. The solid lines in the top panel are fits to  $\delta\nu_a + bP_{\text{mw}}^{0.5}$ . The laser power  $P_L$  was optimized at  $100\ \mu\text{W}$  and  $1.4\ \text{mW}$  for CW and pulsed experiments, respectively. The full-width-at-half-maximum of the laser spot was  $520\ \text{nm}$ , corresponding to  $\sim 2100\ \text{NV}$  centers in the focal volume. The optimized sensitivity  $\eta_{\text{sn}}^{(\text{pulsed})} = 35\ \text{nT Hz}^{-0.5}$  was obtained at  $P_{\text{mw}}$  of  $8.2\ \text{mW}$ , with  $\delta\nu = 718\ \text{kHz}$  and  $C = 0.036$ .

mentioned. The relative improvement in  $\eta_{\text{sn}}^{(\text{cw})}$  over the previous result ( $170\ \text{nT Hz}^{-0.5}$  in Ref. 15) is attributed to the differences in collection efficiency and measurement location within the sample.

As is evident from Eq. (1), the magnetic sensitivity can be improved by optimizing  $\delta\nu$ ,  $C$ , and  $I_0$ . Indeed, the low-field data in Fig. 1 show an interplay between  $\delta\nu$  and  $C$ ; the narrowing of  $\delta\nu$  as decreasing  $P_{\text{mw}}$  is countered by the reduction of  $C$ , and  $\eta_{\text{sn}}^{(\text{cw})}$  takes its minimum at an intermediate value of  $P_{\text{mw}} = 1.64\ \text{mW}$ . We observe  $\delta\nu = \delta\nu_a + bP_{\text{mw}}^{0.5}$  with  $\delta\nu_a \approx 250\ \text{kHz}$  [solid line in the top panel of Fig. 1]. On the other hand, as demonstrated for a single NV center,<sup>24</sup>  $C$  and  $I_0$  can be further improved, respectively, by employing the DNP of  $^{14}\text{N}$  nuclei ( $I = 1$ ) associated with the NV centers and pulsed ODMR.

It is well-established that the excited state of the NV center experiences a level anticrossing near  $50\ \text{mT}$ . In this condition, the optical pumping of the NV electronic spins polarizes the  $^{14}\text{N}$  nuclei into the  $m_I = 1$  state owing to electron-nuclear flip-flops.<sup>25,26</sup> Figure 2(a) plots CW ODMR spectra taken at  $1.5\ \text{mT}$  and  $52.0\ \text{mT}$ , demonstrating clear DNP in the latter. We then repeat the measurements at  $52.0\ \text{mT}$  [ $\Delta$  in Fig. 1] to obtain the minimum sensitivity of  $66\ \text{nT Hz}^{-0.5}$  at  $P_{\text{mw}} = 0.82\ \text{mW}$ .

In the DNP condition, by pumping the three nuclear spin states into a single state, we should ideally achieve a factor-of-three enhancement of  $C$  and thus  $\eta_{\text{sn}}^{(\text{cw})}$ . However, at a given  $P_{\text{mw}}$ ,  $C$  under the DNP is typically only twice as deep as that at low fields. Also, the values of  $C$  giving the minimum sensitivities are  $1.5\%$  at  $1.5\ \text{mT}$  and  $2.6\%$  at  $52.0\ \text{mT}$  [Fig. 2(a)], which are reflected in the obtained sensitivities ( $124\ \text{nT Hz}^{-0.5}$  vs.  $66\ \text{nT Hz}^{-0.5}$ ).  $C$  is determined by a

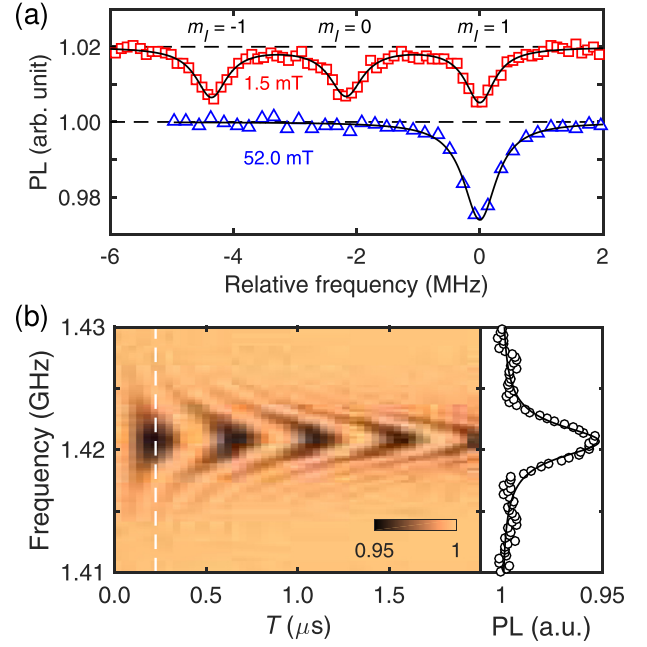


FIG. 2. (a) CW ODMR spectra at  $B_0 = 1.5\ \text{mT}$  and  $P_{\text{mw}} = 1.64\ \text{mW}$  (shifted upward by  $0.02$  for clarity) and at  $B_0 = 52.0\ \text{mT}$  and  $P_{\text{mw}} = 0.82\ \text{mW}$ . (b) An example of pulsed spectroscopy. The cross section at  $T_\pi = 222\ \text{ns}$  is fitted by a Lorentzian.

complicated interplay between  $T_1$ ,  $T_2^*$ , and other optical transition probabilities between the NV electronic energy levels (e.g., Eq. (A5) of Ref. 24). We have examined several physical parameters of our NV ensemble to reproduce the observed  $C$  but have not reached a satisfactory explanation. We leave a detailed analysis on this as a future work.

In CW ODMR at a fixed  $P_{\text{mw}}$ , increasing optical excitation power simultaneously increases  $I_0$  and  $\delta\nu$  while decreasing  $C$ .<sup>24</sup> This leads to an optimal optical power well below the saturation intensity of the NV center. On the other hand, pulsed ODMR temporally separates the optical pumping from the spin manipulation. A higher laser power can be used to significantly increase  $I_0$  while keeping  $C$  and  $\delta\nu$  intact.<sup>24</sup> An example of pulsed spectroscopy,  $C$  as a function of the microwave burst time ( $T$ ) and the microwave frequency, is shown in Fig. 2(b). We denote the experimental sequence as  $\tau_I - T - \tau_R$ , where  $\tau_{I,R}$  are the durations of green laser excitation for spin initialization and readout. By varying the microwave frequency around the  $m_I = 1$  resonance, a chevron pattern typical of pulsed spectroscopy is observed. The cross section at the  $\pi$  pulse condition ( $T_\pi = 222\ \text{ns}$ ) is also shown in Fig. 2(b), from which we deduce  $\delta\nu$  and  $C$  [ $\circ$  in Fig. 1].  $\eta_{\text{sn}}$  for pulsed ODMR is given by<sup>15,24</sup>

$$\eta_{\text{sn}}^{(\text{pulsed})} = \frac{h}{g\mu_B} \frac{\delta\nu}{C} \sqrt{\frac{(\pi\delta\nu)^{-1} + \tau_I + \tau_R}{\tau_R I_0}}, \quad (2)$$

and we obtain the minimum sensitivity of  $35\ \text{nT Hz}^{-0.5}$  for our system.

Through time-resolved fluorescence measurements (data not shown), we optimized  $\tau_I$  and  $\tau_R$  as  $4.5\ \mu\text{s}$  and  $1.5\ \mu\text{s}$ , respectively. These spin initialization and readout times, several times longer than the case of a single NV center

(typically about 1  $\mu\text{s}$  and a few 100 ns, respectively), currently limit the achievable  $\eta_{\text{sn}}^{(\text{pulsed})}$ . The long  $\tau_{\text{I,R}}$  are attributed to the Gaussian profile of the laser spot. Calculations suggest that 58% of the total fluorescence intensity arise from the region outside of the FWHM of the profile, and the NV ensemble existing in this region is subject to substantially lower laser power, resulting in insufficient initialization for shorter  $\tau_{\text{I}}$ . A long initialization time due to the laser spot profile has also been discussed in Ref. 27.

Having demonstrated the potential of this sample for 2D DC magnetic sensing, we next measure  $T_2$  using a Hahn echo sequence ( $\tau_{\text{I}} - T_{\pi/2} - \tau - T_{\pi} - \tau - T_{\pi/2} - \tau_{\text{R}}$ ) and carry out ensemble-based decoherence spectroscopy. Figure 3(a) shows three representative Hahn echo decay curves at around 50 mT, all described well by single-exponential decays. We define  $T_2$  by fitting the echo decay curves as  $A \exp(-2\tau/T_2)$ . A detailed  $B_0$ -dependence of the decoherence rate shown in Fig. 3(c) reveals a multi-peak structure. As depicted in Fig. 3(b), the peak positions [listed in Table I] coincide with the simultaneous spin resonances (SSRs) of the NV and P1 centers:  $S = \frac{1}{2}$  substitutional nitrogen impurities with  $C_{3v}$  symmetry.<sup>28,29</sup> In a dipolarly coupled electron spin system, the refocusing pulse flips resonant spins within its bandwidth, instantly changing local dipolar magnetic fields experienced by the individual NV electronic spins: a process known as the instantaneous diffusion (ID). At the SSR, the NV decoherence is accelerated by

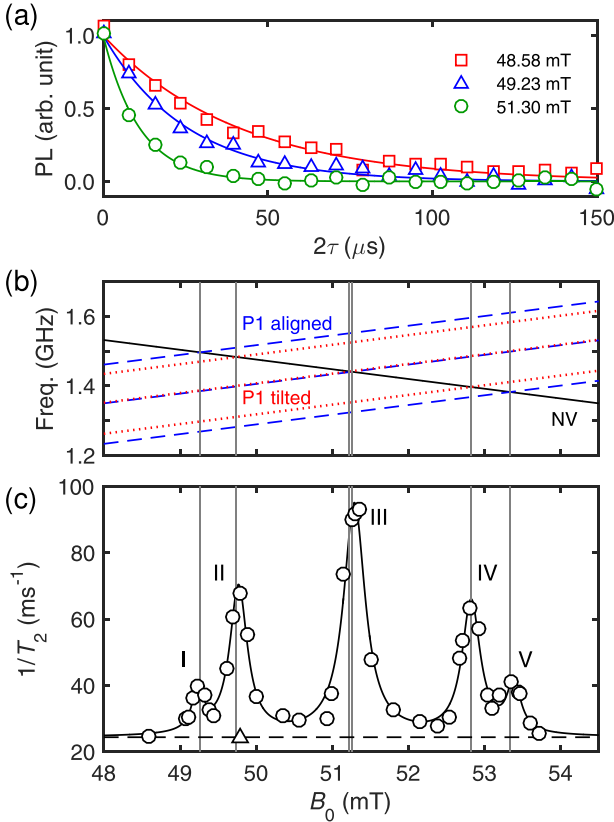


FIG. 3. (a) Hahn echo decays of the NV ensemble at around  $B_0 = 50$  mT. The solid lines are fits by single-exponential decays. (b) The transition frequencies of the NV centers (solid line,  $m_S = 0 \leftrightarrow -1$  with  $m_l = 1$ ) and the P1 centers aligned with  $B_0$  (dashed lines) and tilted from  $B_0$  (dotted lines). The unit and scale of the horizontal axis are the same as (c). The vertical lines indicate  $B_0$  at which the SSRs occur. (c)  $T_2^{-1}$  as a function of  $B_0$ .

TABLE I. Analysis of Fig. 3(c). The peak positions are determined by a quintuple-Lorentzian fit.  $(T_{\text{id}}^{(i)})^{-1}$  specifies  $(T_{\text{id}})^{-1}$  at the  $i$ th peak after subtraction of  $(T_2^{\text{base}})^{-1}$ .  $N_i$  is the impurity density.

Peak	$B_0$ (mT)	$(T_{\text{id}}^{(i)})^{-1}$ ( $\text{ms}^{-1}$ )	$N_i$ ( $10^{17} \text{ cm}^{-3}$ )
I	49.23	12.6	0.19
II	49.77	44.8	0.68
III	51.29	70.0	1.05
IV	52.82	40.2	0.61
V	53.38	14.8	0.22

the increased number of flipped spins. In the case of a homogeneous electron-spin distribution, the ID decay has a form of single-exponential  $\exp(-2\tau/T_{\text{id}})$  with  $T_{\text{id}}$  given by<sup>30,31</sup>

$$\frac{1}{T_{\text{id}}} = DN \sin^2\left(\frac{\beta}{2}\right), \quad \text{with } D = \frac{\pi\mu_0 g^2 \mu_B^2}{9\sqrt{3}\hbar}. \quad (3)$$

Here,  $\mu_0$  is the vacuum permeability,  $N$  is the density of spins rotated, and  $\beta$  is the flip angle of the refocusing pulse. We are thus able to extract the densities of the P1 spins from the increased decoherence rate at the respective peaks.

To do so, we first define the *baseline* decoherence rate  $(T_2^{\text{base}})^{-1}$  as  $T_2^{-1} = 24.4 \text{ ms}^{-1} = (41.1 \mu\text{s})^{-1}$  obtained for the  $m_S = 0 \leftrightarrow 1$  transition at 49.8 mT driven at 4.2658 GHz [ $\Delta$  in Fig. 3(c)]. This spin state shares the same decoherence mechanism with the  $m_S = 0 \leftrightarrow -1$  one, except for the ID due to the P1 spins. The main contributor to  $(T_2^{\text{base}})^{-1}$  is the ID among the NV spins, which, for  $N_{\text{NV}\parallel B_0} = 0.25 \times 10^{17} \text{ cm}^{-3}$  (1/4 of the total NV density<sup>15</sup>), is estimated to be  $20.1 \text{ ms}^{-1}$ . The rest ( $4.3 \text{ ms}^{-1}$ ) should come from (i) the spectral diffusion caused by flip-flops among the P1 spins and (ii) the  $T_1$  relaxation. The lattice  $^{13}\text{C}$  nuclei play a negligible role owing to the isotope enrichment of  $^{12}\text{C}$  in this sample.

We then calculate the P1-induced  $T_{\text{id}}^{-1}$  after subtracting  $(T_2^{\text{base}})^{-1}$  from the measured  $T_2^{-1}$  and the corresponding P1 densities, which are listed in Table I. In Eq. (3),  $\beta = \pi/\sqrt{2}$  is used, due to the smaller rotation angle for the  $S = \frac{1}{2}$  P1 spins relative to the  $S = 1$  NV spins.<sup>31</sup> The ratio  $(N_{\text{II}} + N_{\text{IV}})/(N_{\text{I}} + N_{\text{V}}) = 3.1$  is close to 3: the ratio of the numbers of P1 centers tilted from and aligned with  $B_0$ . On the other hand,  $(N_{\text{I}} + N_{\text{II}} + N_{\text{IV}} + N_{\text{V}})/2 = 0.85 \times 10^{17} \text{ cm}^{-3}$  is less than  $N_{\text{III}} = 1.05 \times 10^{17} \text{ cm}^{-3}$ . The discrepancy may indicate the presence of an additional  $S = \frac{1}{2}$  impurity with the density of  $2 \times 10^{16} \text{ cm}^{-3}$ . We note that, in samples with the P1 density of  $\sim 10^{17} \text{ cm}^{-3}$ , the P1-induced spectral diffusion of the order of a few 100  $\mu\text{s}$  has been observed,<sup>32,33</sup> consistent with our assignment of a-few- $\text{ms}^{-1}$  decoherence rate to this mechanism.

From the total P1 + NV density of  $3.75 \times 10^{17} \text{ cm}^{-3}$ , the dipolar-limited linewidth is estimated to be 95 kHz from the second moment.<sup>34,35</sup> This suggests that the present NV linewidth (250 kHz) is still not limited by the dipolar interactions and there is a room for further reducing the NV linewidth while keeping the NV density at  $10^{17} \text{ cm}^{-3}$ .

Finally, we examine the  $T_{\pi}$ -dependence of  $T_2$ , taking Peak IV as an example.  $T_{\pi}$  of  $\sim 45$  ns used in Fig. 3(c) is so broadband that the linewidth deduced from the multi-Lorentzian fit (16 MHz) does not necessarily reflect the true width. Figure 4(a) demonstrates that, as making  $T_{\pi}$  longer,



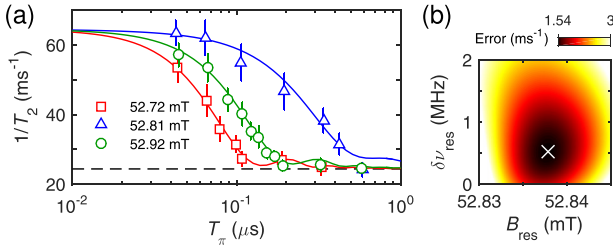


FIG. 4. (a)  $T_\pi$ -dependence of  $T_2^{-1}$  at Peak IV. (b) Fit error as a function of fitting parameters  $B_{\text{res}}$  and  $\delta\nu_{\text{res}}$ . The cross mark ( $\times$ ) indicates the values giving the best fit.

$T_2^{-1}$  gradually falls down to  $(T_2^{\text{base}})^{-1}$ , suggesting that less and less P1 spins are flipped. Considering that an *effective* amount of spins rotated is determined from an overlap between the frequency spectrum of the microwave pulse  $\mathcal{P} = (f_{\text{R}}/f_{\text{R,g}})^2$  and the impurity spin spectrum  $\mathcal{S}$ , we write  $T_\pi$ -dependent  $T_2$  as

$$\frac{1}{T_2} = \frac{1}{T_2^{\text{base}}} + DN_{\text{IV}} \int \mathcal{P}\mathcal{S} \sin^2\left(\frac{\beta}{2}\right) df. \quad (4)$$

Here,  $f_{\text{R}} = (2\sqrt{2}T_\pi)^{-1}$  is the Rabi frequency for  $S = \frac{1}{2}$ ,  $f_{\text{R,g}}(f) = \{(f - \delta f)^2 + f_{\text{R}}^2\}^{1/2}$  with  $\delta f = 2g\mu_{\text{B}}(B_0 - B_{\text{res}})$  is the generalized Rabi frequency,<sup>34</sup>  $\beta(f) = 2\pi f_{\text{R,g}}T_\pi$  is the flip angle, and  $\mathcal{S}(f) = \delta\nu_{\text{res}}/2\pi\{(f - \delta f)^2 + (\delta\nu_{\text{res}}/2)^2\}$  is assumed to be a Lorentzian. Despite the complex form, Eq. (4) contains only two fitting parameters (the resonance magnetic field  $B_{\text{res}}$  and the impurity spin linewidth  $\delta\nu_{\text{res}}$ ) and yet reproduces the experimental data at  $(B_{\text{res}}, \delta\nu_{\text{res}}) = (52.838 \text{ mT}, 520 \text{ kHz})$  [solid lines in Fig. 4(a)]. The extracted P1 linewidth of 520 kHz is broader than the NV linewidth and the dipolar-limited linewidth, hinting at the presence of additional broadening mechanisms in this sample. On the other hand, Fig. 4(b) suggests that, while the error at the best fit is  $1.54 \text{ ms}^{-1}$ , the 200-kHz linewidth can be obtained with the error of  $1.68 \text{ ms}^{-1}$ . Such a small difference in errors can arise, for instance, from the uncertainty in  $T_1$ , which is also  $B_0$ -dependent around the SSR due to cross relaxation.<sup>36–39</sup> For a more refined estimation of the P1 linewidth, fine-tuning of  $B_0$  to the exact P1 resonance and a detailed measurement of  $T_1$  will be helpful. Nonetheless, the method presented here will be a powerful approach to resolve a spin spectrum when applied to external spins.

In summary, by applying both DNP and pulsed ODMR techniques to a near-surface, narrow-resonance-linewidth NV ensemble, we have shown that a photon-shot-noise-limited magnetic sensitivity of  $35 \text{ nT Hz}^{-0.5}$ , highly promising for 2D magnetic imaging, is attainable. We have also measured  $T_2$  and deduced quantitative information on residual paramagnetic impurities in the sample. Decoherence spectroscopy as demonstrated here is applicable to detect magnetic signals external to the sample. Although the present work focused on the internal P1 spins for the purpose of demonstrating the power of decoherence spectroscopy, the magnetic field can be readily tuned to avoid the P1 resonances while still maintaining DNP. Such a condition is suitable to concurrently perform highly sensitive DC magnetic imaging and decoherence spectroscopy of external spins (supplementary material).

See supplementary material for the feasibility of decoherence spectroscopy in detecting external spins.

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